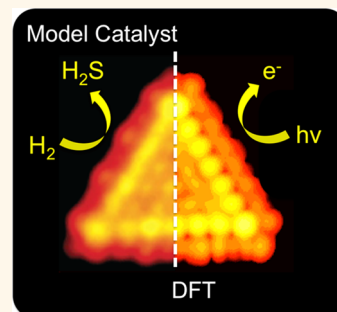


# *In Situ* Detection of Active Edge Sites in Single-Layer MoS<sub>2</sub> Catalysts

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**ABSTRACT** MoS<sub>2</sub> nanoparticles are proven catalysts for processes such as hydrodesulfurization and hydrogen evolution, but unravelling their atomic-scale structure under catalytic working conditions has remained significantly challenging. Ambient pressure X-ray Photoelectron Spectroscopy (AP-XPS) allows us to follow *in situ* the formation of the catalytically relevant MoS<sub>2</sub> edge sites in their active state. The XPS fingerprint is described by independent contributions to the Mo 3d core level spectrum whose relative intensity is sensitive to the thermodynamic conditions. Density Functional Theory (DFT) is used to model the triangular MoS<sub>2</sub> particles on Au(111) and identify the particular sulphidation state of the edge sites. A consistent picture emerges in which the core level shifts for the edge Mo atoms evolve counterintuitively toward higher binding energies when the active edges are reduced. The shift is explained by a surprising alteration in the metallic character of the edge sites, which is a distinct spectroscopic signature of the MoS<sub>2</sub> edges under working conditions.



**KEYWORDS:** molybdenumdisulfide (MoS<sub>2</sub>) · X-ray photoelectron spectroscopy · scanning tunnelling microscopy · DFT · density functional theory · catalysis · nanoparticles · hydrotreating · water splitting

The catalytic chemistry of molybdenite (MoS<sub>2</sub>) is remarkably versatile. MoS<sub>2</sub> has for several decades been in use as a hydrotreating catalyst in oil refineries where it selectively removes both sulfur and nitrogen heteroatoms from the hydrocarbons in crude oil and controls hydrogenation to produce high-value, ultraclean fuels.<sup>1,2</sup> More recently, it has been realized that MoS<sub>2</sub> also performs well as an inexpensive non-noble metal catalyst for the hydrogen evolution reaction (HER) in (photo)-electrochemical cells for water splitting.<sup>3,4</sup> The active sites correspond in these cases to the undercoordinated edge sites of the layers in MoS<sub>2</sub>, which have the capability to dissociate H<sub>2</sub> for hydrogenation reactions, to associate H for hydrogen evolution,<sup>5</sup> and to bind heteroatoms in hydrocarbon molecules during the hydrogenolysis of S and N. In the industrial hydrotreating catalysts, the activity is optimized by increasing the number of edge sites by dispersing single-layer MoS<sub>2</sub> nanoparticles promoted with Co or Ni on a high surface-area  $\gamma$ -alumina support. Stability issues and requirements with respect to conducting properties have led to more elaborate nanostructured MoS<sub>2</sub> materials for use in electrochemistry cells with an

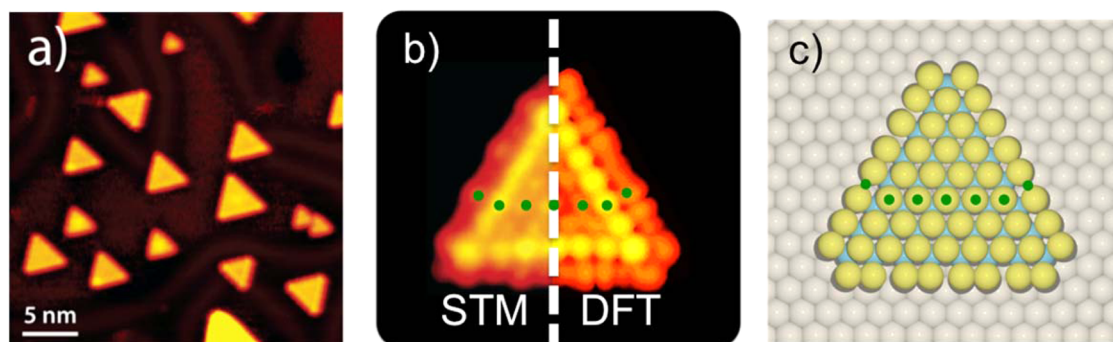
intrinsically high edge concentration<sup>6</sup> or molecular-like MoS<sub>2</sub> clusters consisting of only edge atoms.<sup>7</sup> Fundamental experiments performed on model systems have previously mapped the atomic structure of the edges of single layer MoS<sub>2</sub> nanoparticles in high detail using scanning tunnelling microscopy (STM).<sup>8–11</sup> The STM analysis revealed that the MoS<sub>2</sub> edges are significantly reconstructed compared to bulk truncations of MoS<sub>2</sub> sheets and that the semiconducting nature of MoS<sub>2</sub> is changed to metallic at the edge positions due to electronic edge states.<sup>11,12</sup> These edge states facilitate the adsorption of various molecules relevant to hydrotreating<sup>13–15</sup> and were also recently shown to be the probable cause of nonlinear optical phenomena in MoS<sub>2</sub>.<sup>16</sup> Fourier-Transform Infrared Spectroscopy and Temperature-Programmed Reduction studies have provided valuable chemical information about these edge sites and were used to follow sulfur reduction and hydrogen uptake.<sup>17,18</sup> Furthermore, it was possible to elucidate the sulfur coverage of the same edges in industrial-style model catalysts with High-Resolution Transmission Electron Microscopy (HRTEM)<sup>19,20</sup> again under vacuum conditions. AP-XPS studies performed

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**Figure 1. Model Catalyst and Computational Models:** (a) STM image of MoS<sub>2</sub> nanoparticles synthesized on the Au(111) substrate at 0.15 ML coverage and (b) a close-up of a Mo edge terminated MoS<sub>2</sub> cluster ( $V_{\text{exp}} = -0.52$  V,  $I_{\text{exp}} = 0.44$  nA) compared to the corresponding STM image simulated by means of DFT and the Tersoff-Hamann approximation ( $V_{\text{TH}} = -0.8$  eV,  $I = 17$  pA). The average particle size of  $n = 8$  for a 0.15 ML coverage is reflected, where  $n$  is the number of Mo atoms counted along the edge. The STM simulations fully reproduce all characteristic features seen in the experimental image. This includes the bright metallic brim along the (10 $\bar{1}$ 0) Mo edges, which is a direct reflection of the particular electronic structure at the Mo edges due to metallic edge state,<sup>12</sup> and the off-registry protrusions along the line of basal plane atoms with double periodicity of the S atoms at the edges. The green dots indicate the line of basal plane S atoms and the off-registry protrusion at the edges. The right figure (c) shows the model used for the DFT calculations. Turquoise and yellow spheres correspond to Mo and S atoms, respectively.

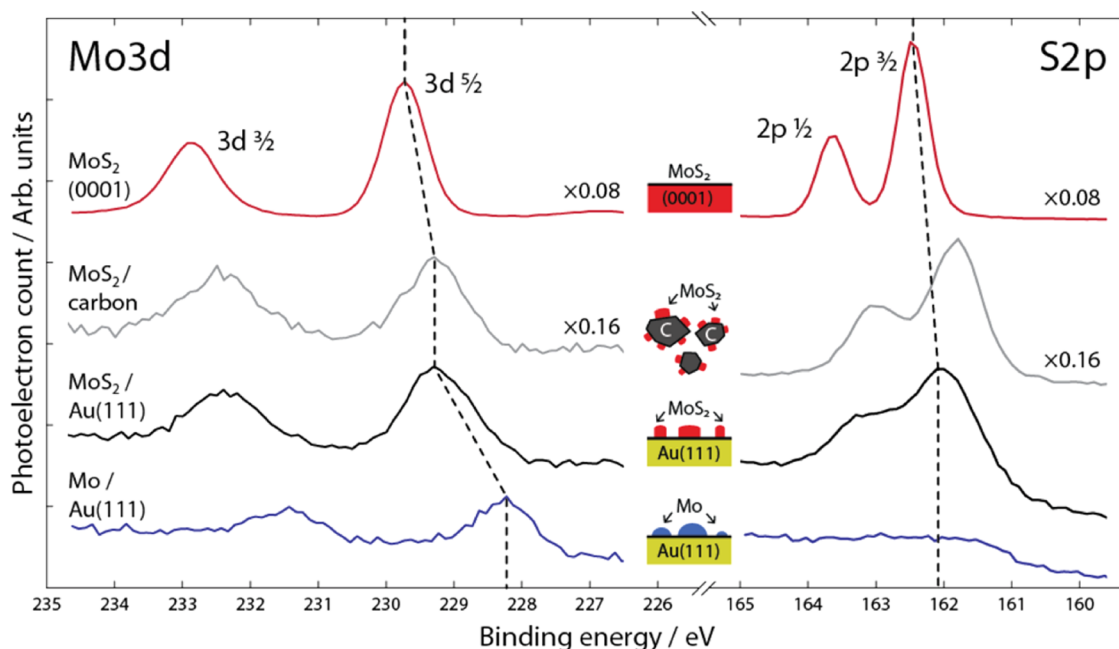
under electrocatalytic hydrogen evolution conditions recently confirmed the presence of MoS<sub>2</sub> as the active catalyst when an amorphous MoS<sub>3</sub> precursor was used,<sup>21</sup> but a spectroscopic signature of the active MoS<sub>2</sub> edges under working conditions is still missing. Here, we use the interplay of AP-XPS, STM and Density Functional Theory (DFT) calculations to elucidate the resulting edge structures which emerge when activating well-characterized single-layer MoS<sub>2</sub> nanoparticles with hydrogen gas at elevated pressures and temperatures that resemble hydrotreating catalysis conditions and the reductive environment under hydrogen evolution conditions. We first show that the Mo 3d XPS spectra can be described by independent contributions that depend on the position of the Mo atoms within the single-layer MoS<sub>2</sub> nanoparticles, where edge Mo atoms exhibit chemical shifts distinct from those in the basal plane. We then identify how such shifts evolve upon reduction of the nanoparticle edges, which enables us to monitor the activation of our MoS<sub>2</sub> model catalyst using ambient-pressure XPS. We demonstrate that the fully sulfided starting structure and the active edge structures with a lower S coverage can be discriminated on the basis of XPS data.

## RESULTS AND DISCUSSION

The large-scale STM image in Figure 1a illustrates the morphology of the MoS<sub>2</sub> nanoparticles used in this study. This MoS<sub>2</sub> model catalyst system consists of (0001)-oriented single-layer MoS<sub>2</sub> nanoparticles (*i.e.*, a S–Mo–S layer) on a Au(111) surface and can be synthesized by evaporation of Mo and annealing at 673 K in a 10<sup>−6</sup> mbar H<sub>2</sub>S atmosphere. The advantage of using this model system is that it consists of very well-defined MoS<sub>2</sub> nanoparticles, the edge structure of which has previously been characterized in atomic detail.<sup>11,12</sup> The MoS<sub>2</sub> nanoparticles have a very strong

preference for a triangular shape and are predominantly terminated by the so-called (10 $\bar{1}$ 0) Mo-edges. The Mo-edges are also established in (S)TEM studies to be by far the most predominant edge type in industrial-type MoS<sub>2</sub> catalysts, so our model system exposes the catalytically most relevant sites.<sup>19,20</sup> The MoS<sub>2</sub> morphology is in fact slightly sensitive to the size of the nanoparticles,<sup>9</sup> but for the coverage of  $\sim 0.15$  ML illustrated in Figure 1a, approximately 95% of the edge sites expose the (10 $\bar{1}$ 0) Mo-edges. In the as-synthesized and fully sulfided state seen in the atom-resolved STM image in Figure 1b, the Mo-edge is fully covered with sulfur (100% S) binding S<sub>2</sub> dimers at each Mo edge atom (Figure 1c). DFT calculations of realistic computational models such as the one illustrated in Figure 1c allow simulating STM images and obtaining core-level Binding Energy (BE) shifts that can be directly compared to the model catalyst results. These novel representative models simultaneously consider the finite dimension of MoS<sub>2</sub> nanoparticles and the presence of the Au(111) support. We have systematically calculated the electronic structure of Au(111)-supported MoS<sub>2</sub> nanoparticles of different edge composition representative of the fully sulfided edges (100% S) prepared experimentally (Figure 1a) as well as Mo-edges predicted to be present under H<sub>2</sub> reducing conditions.<sup>8,22</sup> For the as-synthesized MoS<sub>2</sub> nanoparticle in Figure 1b, our STM simulations reproduce all characteristic features seen in the experimental image. Such excellent agreement indicates that the theoretical models are fully representative of the experimental MoS<sub>2</sub> nanoparticles and can thus be used to perform the interpretation of XPS results.

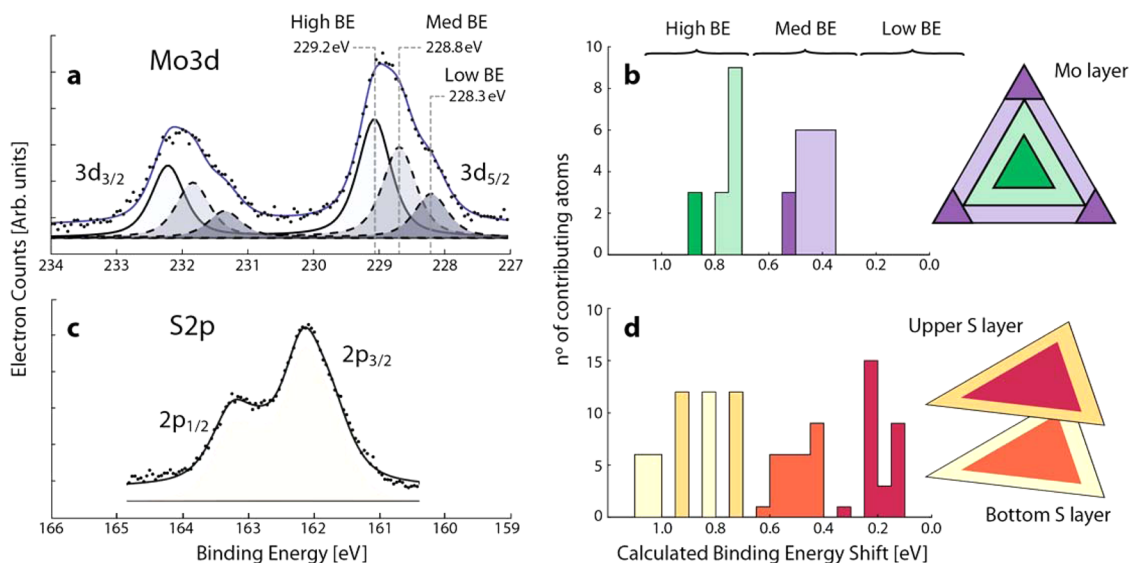
**X-ray Photoemission Spectroscopy at Ultrahigh Vacuum Conditions.** First, we recorded a comparative series of XPS spectra of a cleaved MoS<sub>2</sub>(0001) single crystal surface as a reference, industrial-style MoS<sub>2</sub> catalyst nanoparticles



**Figure 2.** XPS core level spectra for different MoS<sub>2</sub> species: XPS core level spectra of the Mo 3d peaks (3d<sub>3/2</sub> and 3d<sub>5/2</sub>) and S 2p peaks (2p<sub>1/2</sub> and 2p<sub>3/2</sub>) for a MoS<sub>2</sub>(0001) surface, MoS<sub>2</sub> nanoparticles on a C support, and metallic Mo and MoS<sub>2</sub> on Au(111), respectively. The spectra have been measured at the SX700 beamline at Aarhus Storage Ring Denmark (ASTRID). Photon energies were  $h\nu = 340$  eV for Mo 3d and  $h\nu = 300$  eV for S 2p. All spectra were recorded in UHV at 298 K.

on a carbon support, the MoS<sub>2</sub>/Au system and metallic Mo/Au, respectively. The doublet structure of the Mo 3d and S 2p spectra illustrated in Figure 2 reflects the characteristic spin–orbit splitting between Mo 3d<sub>3/2</sub> and 3d<sub>5/2</sub> states and S 2p<sub>1/2</sub> and 2p<sub>3/2</sub>, respectively. The peaks of the reference MoS<sub>2</sub>(0001) surface were sharply resolved and symmetric with positions recorded at 229.6 eV for Mo 3d<sub>5/2</sub> and at 162.5 eV for S 2p<sub>3/2</sub>, in excellent agreement with literature values for MoS<sub>2</sub>(0001).<sup>23</sup> For the MoS<sub>2</sub>/Au samples, the Mo 3d<sub>5/2</sub> main peak lies  $\sim 0.4$  eV lower in BE at 229.2 eV and the approximately same peak shift is also seen in the S 2p signals. This is mainly due to the screening effects arising from the interaction with a metallic substrate and is also an indication of partial charge transfer from the substrate which we observe also in the DFT calculations of MoS<sub>2</sub> on Au presented here; upon interaction with the completely sulfided MoS<sub>2</sub> nanoparticle, some electron density is depleted from Au and accumulated mainly in the Au–S bonding region. Such charge transfer is expected from XPS experiments from the inverse system (gold on MoS<sub>2</sub>)<sup>24,25</sup> and also from the DFT-based studies by Tsai *et al.* involving stripe models of MoS<sub>2</sub> on Au(111).<sup>26</sup> In addition, the Mo 3d peaks present a distinct shift and asymmetry toward lower binding energies compared with the MoS<sub>2</sub>(0001) reference for the MoS<sub>2</sub>/Au. A slighter asymmetry is also seen for the MoS<sub>2</sub>/C technical catalyst system, which also contains nanoparticulate MoS<sub>2</sub>, but with a larger average size.<sup>19</sup> Furthermore, whereas the S 2p doublet is well-resolved on the MoS<sub>2</sub>(0001) sample, the 2p<sup>1/2</sup> and 2p<sup>3/2</sup> peaks broaden and overlap in MoS<sub>2</sub>/C and

even more so in MoS<sub>2</sub>/Au—an indication that S is present in different chemical environments in these samples. To elucidate the origin of the asymmetrical peak shape shown in Figure 2, we conducted a high-resolution XPS study for the MoS<sub>2</sub>/Au system at the Advanced Light Source (Berkeley) using the same synthesis procedure. We start by analyzing the spectra of the as-synthesized sample (Figure 3a), where particles with completely sulfided Mo-edges (100% S) dominate. The peak fitting (see Methods) reveals that the asymmetrical Mo 3d<sub>5/2</sub> peak is best fitted by three components, *i.e.* two main components at 229.2 and 228.8 eV, which are present with a  $\sim 1:1$  ratio, and a smaller peak at 228.3 eV. The core-level shifts (CLS) calculated with DFT (see methods) are used in order to assign these three peaks, which will henceforth be referred to as high, medium, and low BE components. The calculated Mo 3d CLS within the MoS<sub>2</sub> nanoparticles with 100% S edges are consistently lower (*i.e.*, they are more metallic) for Mo atoms in edge positions than for those in the basal plane (Figure 3b). Such differences divide the calculated spectra into two regions which are separated by  $\sim 0.4$  eV, in agreement with the separation between the high and medium peak components fitted from the experimental spectra. Therefore, for the as-synthesized MoS<sub>2</sub> sample, the high BE component, which is closest to the MoS<sub>2</sub>(0001) reference, is assigned to basal plane Mo atoms, whereas the medium BE peak is attributed to edge Mo atoms. Previous results are in line with this assignment, where one or more low BE components arise when a bulk MoS<sub>2</sub>(0001) sample is crushed, leading to larger abundance

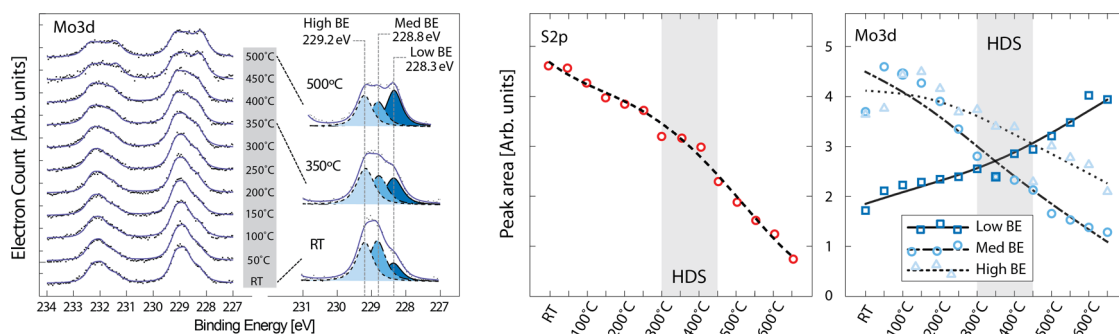


**Figure 3.** Deconvolution of the XPS spectra and DFT-calculated shifts: Mo 3d and S 2p XPS spectra (a and c) recorded at room temperature and DFT-calculated core-level binding energy (BE) shifts (b and d) for Au(111)-supported MoS<sub>2</sub> nanoparticles. The Mo 3d fitted spectrum is deconvoluted into three different components (high, medium, and low BE), which are assigned according to the calculated shifts. The different colors within the histograms in the right panels denote the location of Mo and S atoms (corner, edge, near-edge, or basal plane) in the MoS<sub>2</sub> triangle contributing to the different Mo 3d and S 2p shifts, according to the color-coding indicated within the corresponding triangle drawings.

of undercoordinated edges.<sup>23</sup> However, it should be noted that, as is shown in the forthcoming discussion involving the AP-XPS experiments, these apparently straightforward assignments are subject to variation upon changes in edge composition (sulphidation state) that modify the BE shifts. Importantly, the origin of the 0.4 eV shift to lower BE of the edge Mo atoms can be traced back to their metallic character. This is also directly reflected by the bright brim at the MoS<sub>2</sub> edge in the atom-resolved STM image (Figure 1b), which is caused by one-dimensional metallic edge states specific to the 100% Mo-edge. In the projected density of states (pDOS) of a MoS<sub>2</sub> nanoparticle, a distinctively higher density across the Fermi level is found for the d orbitals of Mo atoms at the edges (Supporting Information Figure S1), but not for those at basal plane positions. The effect resembles the so-called surface core level shift, which is observed in XPS spectra on transition metal surfaces<sup>27</sup> and is a powerful fingerprint for monitoring chemistry on such surfaces.<sup>28,29</sup> In MoS<sub>2</sub>, the metallic character of the edges is also directly expressed in the CLS of the Mo atoms, which are thus directly discernible in high-resolution XPS. This enables the use of *in situ* XPS methods to monitor the nature of the active sites within MoS<sub>2</sub>-based catalysts. We note that the absolute shifts for the two Mo sites (edge and basal plane) with respect to metallic bulk Mo are underestimated in the calculations by 0.2 eV. Such underestimation could originate from the reference model chosen for metallic Mo (an Au-supported Mo<sub>21</sub> nanoparticle), which might not be fully representative of the experimental Mo peak. In addition, the absolute position of the core levels is also influenced by the

screening effects leading to the recently observed band structure renormalization for single-layer MoS<sub>2</sub>(0001).<sup>30</sup> However, we expect this effect to be similar for edge and basal plane Mo atoms and it should therefore not significantly affect the relative shifts between different Mo sites.

In the S 2p spectra of MoS<sub>2</sub>/Au in Figure 2 and the corresponding high resolution spectrum in Figure 3c, we observe a significant broadening compared to the well-resolved S 2p doublet of the MoS<sub>2</sub>(0001) reference, which is indicative of S in several different chemical environments.<sup>31</sup> Our calculated CLS values support this observation, showing a rather wide distribution of S 2p shifts within a 0.8 eV range depending on the position of S in the MoS<sub>2</sub> particle (Figure 3d). Unlike for amorphous MoS<sub>3</sub> and MoS<sub>2</sub>,<sup>21,31</sup> which can be distinguished on the basis of their S 2p spectra,<sup>21,30</sup> the scattering of spectroscopically different S species within Au-supported MoS<sub>2</sub> nanoparticles is too large to allow a straightforward interpretation of the spectral shape. Nevertheless, trends are that S in the S<sub>2</sub> dimers on the edges are placed 0.1–0.5 eV higher in BE relative to the basal plane counterparts. Similarly, broad spectra were calculated for the S species at the Mo-edges of MoS<sub>2</sub> stripe models, where S 2p BEs were found to be larger for the S atoms forming dimers than for monomeric edge S.<sup>21</sup> Another effect in our calculated S 2p CLS is that the interaction with the Au support splits the basal plane S atoms into a higher BE component (for the bottom layer S atoms in contact with Au(111)) and a lower BE component (for the upper S layer). The shifts toward larger BE for those S atoms interacting with Au is a result of the formation of S–Au bonds, which drives



**Figure 4.** *In-situ* monitoring of the edge configuration through AP-XPS spectra: Peak fits of the Mo3d spectra between RT and 500 °C (left), together with plot of the Mo3d and S2p peak areas determined from the peak fits (right), in the range from RT to 650 °C. A close-up of the deconvoluted Mo3d<sub>5/2</sub> spectra for three characteristic temperatures (Room Temperature, 350 °C, and 500 °C) is shown. These plot close-ups clearly show the differences in the relative intensities of the medium and high BE shifts, which are due to the evolution towards the high BE region experienced by the edge Mo atoms upon reduction from 100% S edge composition to 50% S. The shaded area denoted “HDS” indicates the typical temperatures used in industrial hydrotreating catalysis.

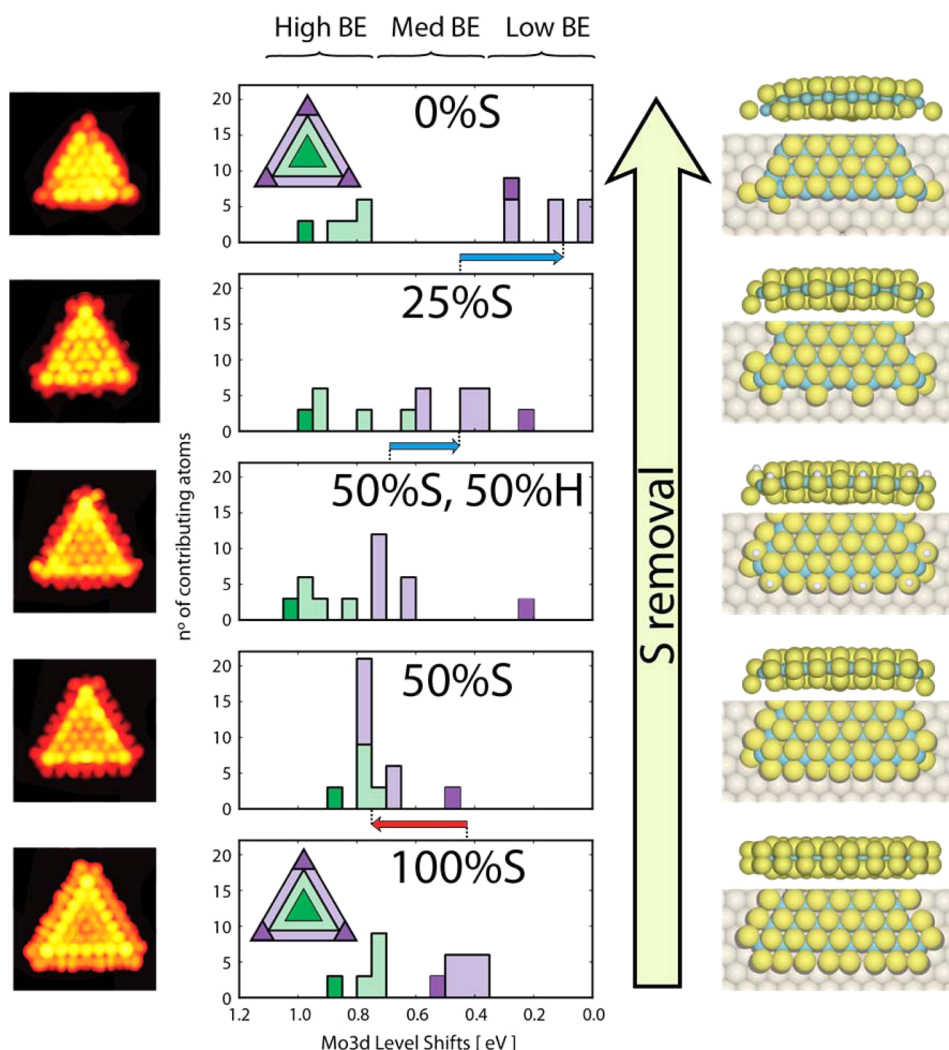
charge away from Au to the S–Au bonding region and polarizes the valence  $p_z$  orbitals of the lower S layer toward the Au surface.

#### X-ray Photoemission Spectroscopy at Elevated H<sub>2</sub> Pressure.

To follow the activation of the catalyst by exposure to hydrogen gas, we performed AP-XPS on the MoS<sub>2</sub>/Au sample by sequentially heating it in 0.01 Torr H<sub>2</sub> and recording the Mo 3d and S 2p XPS spectra for every 50 °C from room temperature to 650 °C. For these thermodynamic conditions, it is predicted that edge S will be titrated away as H<sub>2</sub>S and some H atoms may adsorb on remaining edge S atoms.<sup>8,18,32</sup> Therefore, we directly compare the AP-XPS observations with the calculated shifts in BE for the same Au(111)-supported MoS<sub>2</sub> nanoparticle models but with different edge compositions. In particular, we have considered Mo-edges with a sequentially lowered S coverage and edges with S–H groups corresponding to 100% S, 50% S, 50% H to 50% S, 25% S, and 0% S (see optimized structures in Figure 5). It should be noted that under UHV conditions ( $P \sim 10^{-10}$  mbar), STM experiments show that coalescence and decomposition of the MoS<sub>2</sub> nanoparticles is observed only for temperatures over 550 °C. In contrast, under the elevated H<sub>2</sub> pressures used here, the most pronounced changes in the XPS spectra occur in the catalytically interesting temperature range between 200 and 450 °C (Figure 4), where the MoS<sub>2</sub> nanoparticle morphology and edge structure is otherwise unaffected without the hydrogen gas. The observed changes in the XPS data in Figure 4 do not seem to require any additional peaks for the deconvolution of the Mo 3d spectra. Thus, the spectra have been deconvoluted to the same three peaks from the fit at room temperature and UHV conditions (Figure 2). The positions of these peaks are kept constant for the fitting throughout the HP-XPS series, and therefore, only their intensities change. This means that Mo 3d shifts within the supported MoS<sub>2</sub> NP are registered only as variations in the relative intensities between the different components (*i.e.*, high, medium, or low BE). Indeed, we

find that the ratios between the fitted Mo 3d peaks, as well as the total intensity of the S 2p spectra, change systematically from the room temperature situation. In particular, the Mo 3d spectral shape in Figure 4 clearly changes its weight from high to lower BE. In turn, the corresponding S 2p spectra (Supporting Information Figure S2) remain broad and without apparent shape changes that would require the consideration of additional peaks. Nevertheless, the total intensity is gradually reduced, indicating that S is indeed being removed from the MoS<sub>2</sub> nanoparticles.

A closer look at the qualitative changes in the Mo 3d spectral shapes reveals that exposure to H<sub>2</sub> at increasing temperatures initially leads to a decrease in intensity for the medium BE component. This behavior can be explained as a change in edge coverage from 100% S (dimers) toward nanoparticles with 50% S (monomers). The DFT calculations reveal a shift toward higher Mo 3d BEs for the Mo atoms at the edges when these are reduced from 100% S to 50% S (Figure 5, red arrow). In the *in situ* XPS spectra, this appears as a change in the relative intensity of the high and medium BE components. Reduction processes normally lead to the opposite effect in the core level positions, *i.e.*, shifts toward more metallic character, but in this case, such an unexpected shift toward higher BE has its origin in the partial loss of the one-dimensional metallic state that is characteristic of fully sulfided (10 $\bar{1}$ 0) Mo-edges. This is clearly manifested in the simulated STM images, which show that, upon reduction from 100% S to 50% S edge composition, the bright brim along the edge of the MoS<sub>2</sub> nanoparticles practically disappears, as the edges become semiconducting (Figure 5). The transition from metallic to nonmetallic character of the edges is due to the substitution of the edge S<sub>2</sub><sup>δ-</sup> dimers by S<sup>2-</sup> anions, which have a stronger oxidizing effect on the neighboring Mo atoms. This is also evidenced in the calculated charge density around the different edge S species; the Bader charge of the S<sup>2-</sup> within the 50% S edge (−0.52e) is significantly



**Figure 5.** Simulated STM spectra and core-level shifts: Evolution of the calculated Mo 3d binding energy shifts for different positions within  $n = 8$  MoS<sub>2</sub> nanoparticles of varying edge composition. From bottom to top, models with edge compositions of 100% S, 50% S, 50% S to 50% H, 25% S, and 0% S have been considered. The different colors within the histograms in the right panels denote the location of Mo atoms (corner, edge, near-edge, or basal plane) in the MoS<sub>2</sub> NP contributing to the different Mo 3d shifts, according to the color-coding indicated in triangle drawings shown in the top and bottom graphs. To the left are shown STM images calculated with the Tersoff-Hamann approach at  $I = 17$  pA and  $V = -0.8$  eV. The red and blue arrows indicate the direction of the changes in BE shift for the edge sites during S removal. Turquoise, yellow, and white spheres correspond to Mo, S, and H atoms, respectively.

more negative than around the S dimers in the 100% S ( $-0.32$ e). The described transition is further confirmed by the pDOS on the valence d orbitals of the edge Mo atoms (Supporting Information Figure S1). For the 100% S edge, there is a significantly higher pDOS across the Fermi level for Mo atoms at the edges than for those on basal plane positions, whereas for the 50% S particle, the pDOS at the Fermi level is similarly low for Mo atoms in either location. The partial hydrogenation of the 50% S edge which MoS<sub>2</sub> particles undergo when exposed to H<sub>2</sub> also affects the CLS. The presence of S–H groups mostly perturb the *near-edge* Mo atoms, whose peaks are shifted toward even higher BEs (Figure 5), leading to a situation where the edge Mo are again metallic relative to the basal plane Mo. This is consistent with the presence of metallic edge

states for this edge configuration<sup>33</sup> and also reflected by the higher pDOS across the Fermi level (Supporting Information Figure S1). We have also calculated the S 2p shifts for the 50% S and 50% S to 50% H models, and compared them to those of the completely sulfided 100% S edge results. The calculated S 2p shifts for the different models (presented in Figure S3 of the Supporting Information) reveal that the BEs for the basal plane S atoms are not significantly affected by the edge reduction. In contrast, the S 2p of the S<sup>2-</sup> anions of the edges of the 50% S model are shifted  $\sim 2$  eV shift toward lower binding energies with respect to the S<sub>2</sub><sup>δ-</sup> dimers of the 100% S edge, which is in agreement with the shifts calculated for the same species in unsupported MoS<sub>2</sub> stripes.<sup>21</sup> The hydrogenation of the edges giving rise to the 50% S to 50% H model leads to the

opposite effect, with shifts toward higher BE that merge the S 2p states of edge and basal plane S atoms in the same energy region. Additional peaks at low BE do not appear in the experimental S 2p HP-XPS spectra (Supporting Information Figure S2) upon desulfurization, thereby suggesting that the edges of the desulfurized MoS<sub>2</sub> nanoparticles become partially hydrogenated at elevated temperature in the hydrogen gas (*i.e.*, represented by the 50% H to 50% S configuration).

The AP-XPS Mo 3d spectra at further increased temperatures are characterized by a high intensity of the low BE component. At these temperature regimes (beyond 450 °C), highly desulfurized nanoparticle compositions (*e.g.*, 25% S, 0% S and further decomposed species) are expected to dominate, which in the DFT calculations give rise to more metallic shifts as the S concentration is reduced (Figure 5, blue arrows). In concordance, these lower sulfided compositions also give rise to higher pDOS on the valence d states across the Fermi level for Mo atoms at the edge positions (Supporting Information Figure S1) and to brighter features in the simulated STM images shown in Figure 5.

The peak area plots in Figure 4 show the quantitative area changes of the Mo 3d peak fit components and the S 2p spectral intensity as a function of temperature. The values for the S 2p peak areas presented in Figure 4 correspond to the total integrated area below each spectrum of Supporting Information Figure S2, which is related to the total amount of S in the model catalyst. The main observation for the Mo 3d spectra is that the medium BE (circles) and high BE (triangles) components decrease in intensity, while the intensity of the low energy component (squares) grows. The initial decrease of the two higher energy peaks is slow, but from 200 °C and onward up to 450 °C, the slope of the middle BE becomes much steeper and this component drops by almost 50%. At the same time, the low BE signal increases, indicating that some of the particles are reduced to less sulfided states, in particular above 300 °C. Overall, this is consistent with the conversion of the middle BE component associated with Mo-edges with 100% S into edges with 50% S coverage or lower within the catalytically interesting temperature interval. We do not consider a sintering process to be the likely explanation for the changes in basal plane to edge ratio, since STM investigations have shown that the initial size of the MoS<sub>2</sub> triangles are unaffected by extensive hydrogen exposure (hours) at 10<sup>-6</sup> mbar and 450 °C. It should be noted

that the overall process also results in a drop of the total Mo 3d spectral area, which likely reflects a loss of Mo due to accelerated alloying between the gold substrate and the molybdenum atoms exposed by the reducing conditions.<sup>34</sup> As a result of the intense desulfurization which occurred in the experiment without a balancing source of S (*e.g.*, S-containing hydrocarbons as in hydrotreating), the MoS<sub>2</sub> nanoparticles may even have decomposed beyond the reduced edges seen in Figure 5 at the highest temperatures (>500 °C). In fact, the S 2p/Mo 3d ratio decreases by approximately a factor 2.2 from the first to the last measurement in the series. This change cannot be explained by removal of only edge S in a representative *n* = 8 MoS<sub>2</sub> triangle (Figure 1), and at high temperatures in our experiment, the reaction must have proceeded further than this, thereby desulfurizing the near-edge region of the basal plane.

## CONCLUSIONS

We have shown that the detailed spectral shape of Mo 3d core level results from the nanoscale structure of MoS<sub>2</sub>. Moreover, high-resolution XPS shows that the spectral shape is sensitive to the S coordination at the edges, and this effect enables characterization of the transitions involved in desorption of S and adsorption of H that describes the catalytically active MoS<sub>2</sub> edges. Remarkably, the edge reduction process from 100% to 50% sulfur coverage leads to a shift toward higher Mo 3d BE, which is opposite to the conventional interpretation of "oxidation" in CLS. The shift can be traced back to variations in the metallic character of the Mo edge, which is lost when S<sub>2</sub><sup>δ-</sup> dimers are replaced by more reducing S<sup>2-</sup> species. We anticipate that the chemical shifts seen in the molybdenum core levels by XPS should also be observable in Near-edge X-ray absorption (NEXAFS) or Resonant Inelastic X-ray Scattering (RIXS) experiments, which are already powerful tools for the characterization of catalysts in the working state.<sup>35</sup> These techniques also probe the electron levels near the Fermi level, and should therefore be directly sensitive to the changes in the metallic character of the edges during formation of the active sites. They are furthermore applicable for characterization of technical catalysts beyond the planar model catalysts approach applied in the present studies, and since the observed change in the metallic character is predicted to be independent of the Au support,<sup>33</sup> the present work thus provides a general framework for interpretation of *in situ* X-ray spectroscopy of MoS<sub>2</sub> edges.

## METHODS

**MoS<sub>2</sub>/Au(111) Model System.** The STM experiments were performed in an ultrahigh vacuum (UHV) chamber equipped with a home-built Aarhus STM. The XPS measurements were measured

at two different beamlines, the SX700 beamline at ASTRID (Aarhus Storage Ring, Denmark) and the 11.0.2 beamline at the ALS (Advanced Light Source, Lawrence Berkeley National Laboratory). For direct comparison and to avoid contamination problems due to air exposure, the MoS<sub>2</sub> nanoparticles were

synthesized *in situ* for each experiment in either of the UHV chambers using a precalibrated Mo e-beam evaporator (Oxford Research EGN4 or EGCO4), a homemade H<sub>2</sub>S doser<sup>36</sup> on the same freshly prepared Au(111) sample. The synthesis is performed in two steps, first by e-beam evaporation of Mo in a sulphiding atmosphere (H<sub>2</sub>S,  $P = 10^{-6}$  mbar) followed by continued sulfidation at elevated temperature (450 °C) (see details in refs 9 and 11). We synthesized ~0.15 ML of MoS<sub>2</sub> on the gold substrate, where 1 ML is equivalent to a fully covered substrate. Regular checks were made for drift in the Mo flux with the flux reading of the e-beam evaporator and a quartz crystal microbalance (ALS).

**X-ray Photoelectron Spectroscopy.** We have recorded spectra of the Mo 3d peak with a photon energy of 370 eV and the S 2p peak with 300 eV. For each of the Mo 3d and S 2p spectra, a Au 4f spectrum was recorded with the same photon energy to allow energy calibration. All XPS peaks were fitted with pseudo-Voigt line-shapes, consisting of a sum of a Gaussian and Lorentzian contribution, using the same method as implemented by others.<sup>37</sup> For the MoS<sub>2</sub> nanoparticles, we use symmetric peaks while asymmetric peaks were used for metallic molybdenum, consistent with the Fermi sea screening of the core hole due to the coupling to the gold substrate's free electrons.<sup>38</sup> A linear background was subtracted from all spectra. The splitting between the Mo 3d<sub>3/2</sub> and 3d<sub>5/2</sub> peaks was determined to be 3.16 eV. This is consistent with the value reported in literature.<sup>23,39</sup> For each spin-orbit doublet, the ratio between the two peaks was found narrowly around the ratio given from quantum numbers, *e.g.*, 2:3 for the Mo 3d peak. The Au 4f spectra were fitted with two asymmetric components, where the highest binding energy corresponding to the bulk component was assigned a binding energy of 84.0 eV to which the intensity and energy of the rest of the spectra were calibrated.

**DFT Calculations.** The DFT calculations were performed using the PBE exchange-correlation functional, the Projector Augmented Wave (PAW) method,<sup>40</sup> and a real space grid for the expansion of the wave functions, as implemented in the GPAW code.<sup>41</sup> To calculate the relative binding energy of the Mo 3d and S 2p electrons for different Mo and S atomic species, a PAW-type pseudopotential with an electron hole in the selected level was used for describing the atom the core electron was removed from. A compensating homogeneous background density was included in order to compensate the positive charge resulting from the removal of an electron from a core level. Prior to the generation of the core-holes, the geometries for the different models used were optimized until the forces on each atom were lower than 0.025 eV/Å. The upper layer of the two-layer slab used to represent the Au(111) substrate was also allowed to relax. For each optimized system, the generation of an electron hole was considered for all nonequivalent atomic positions. The chemical shifts of Mo 3d were calculated relative to the average shift of Mo atoms within an Au(111)-supported Mo<sub>21</sub> cluster, which corresponds to metallic Mo. A spectator Mo atom located at the opposite side of the slab of each system was used to calculate the relative shifts between each MoS<sub>x</sub> nanostructure and the metallic Mo<sub>21</sub> cluster. The shifts of S 2p were calculated relative to a spectator S atom on the opposite side of the slab. The STM simulations in Figures 1 and 5 are based on the Tersoff-Hamann approximation, which relates the intensities of STM images with isosurfaces of the local density of states at the Fermi level.<sup>42,43</sup>

**Conflict of Interest:** The authors declare no competing financial interest.

**Supporting Information Available:** pDOS plots, HP-XPS S 2p spectra, and calculated S 2p shifts are presented in Figures S1, S2, and S3, respectively. The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsnano.5b03199.

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